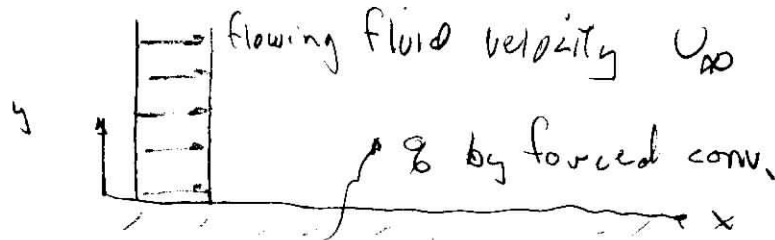
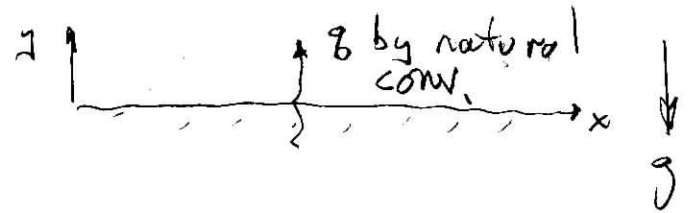
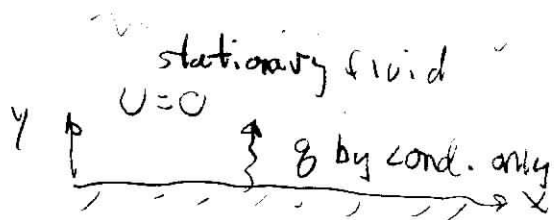


Convection

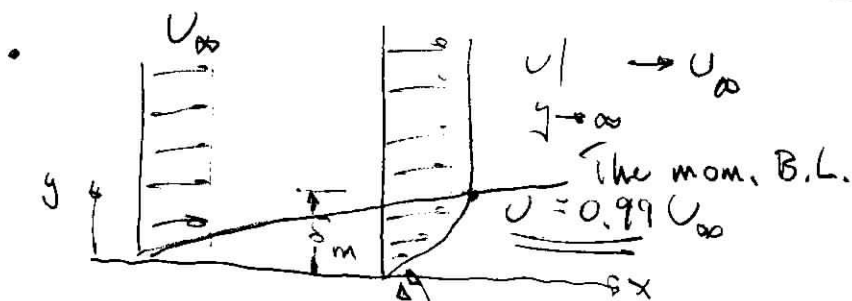
fluid rise



Fluid viscosity (absolute) μ (Dynamic) $\left[\frac{\text{kg}}{\text{m s}} \sim \frac{\text{N s}}{\text{m}^2} \sim \text{Pa s}, 1 \text{ Pa s} = 10 \text{ poise} \right]$

Kinematic $\nu = \frac{\mu}{\rho}$ $\left[\frac{\text{m}^2}{\text{s}} \sim \text{stoke}, 1 \text{ stoke} = 1 \frac{\text{cm}^2}{\text{s}} = 0.0001 \frac{\text{m}^2}{\text{s}} \right]$

Wall shear stress $\tau_{\text{wall}} \equiv \mu \left. \frac{\partial U}{\partial y} \right|_{y=0}$ $[\text{N/m}^2]$
 (Newtonian fluids $\sim$ linear relation)



note $U = U(y)$ only
 $V = 0$
 $w = 0$

$\underline{V} = U \hat{i} + V \hat{j} + w \hat{k}$

$U=0$, no slip condition!
 τ is prop. to $\left. \frac{\partial U}{\partial y} \right|_{y=0}$ here

But you need to know $U(y)$ to calc. τ_{wall} .

or

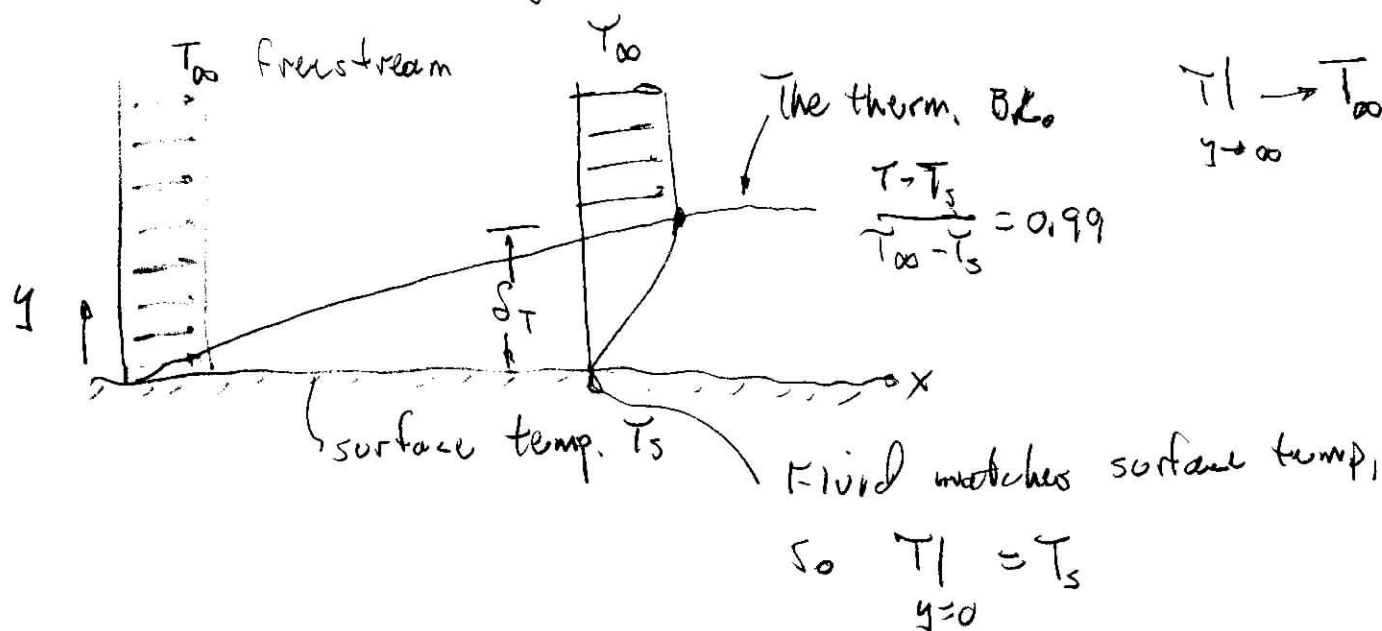
2/

$V \sim$ upstream velocity
(far to left)

$$F_{\text{friction}} = \frac{1}{2} C_f A_s 3V^2 \quad [N]$$

Ultimately Γ_f will be related to a heat transfer coeff

Thermal Boundary Layer



Prandtl ~~the~~

$$Pr \sim \frac{\text{mom. diffusion}}{\text{therm. diffusion}} = \frac{\nu}{\alpha} = \frac{\mu C_p}{k} \quad \text{dimensionless}$$

Lig metals

gases, 221

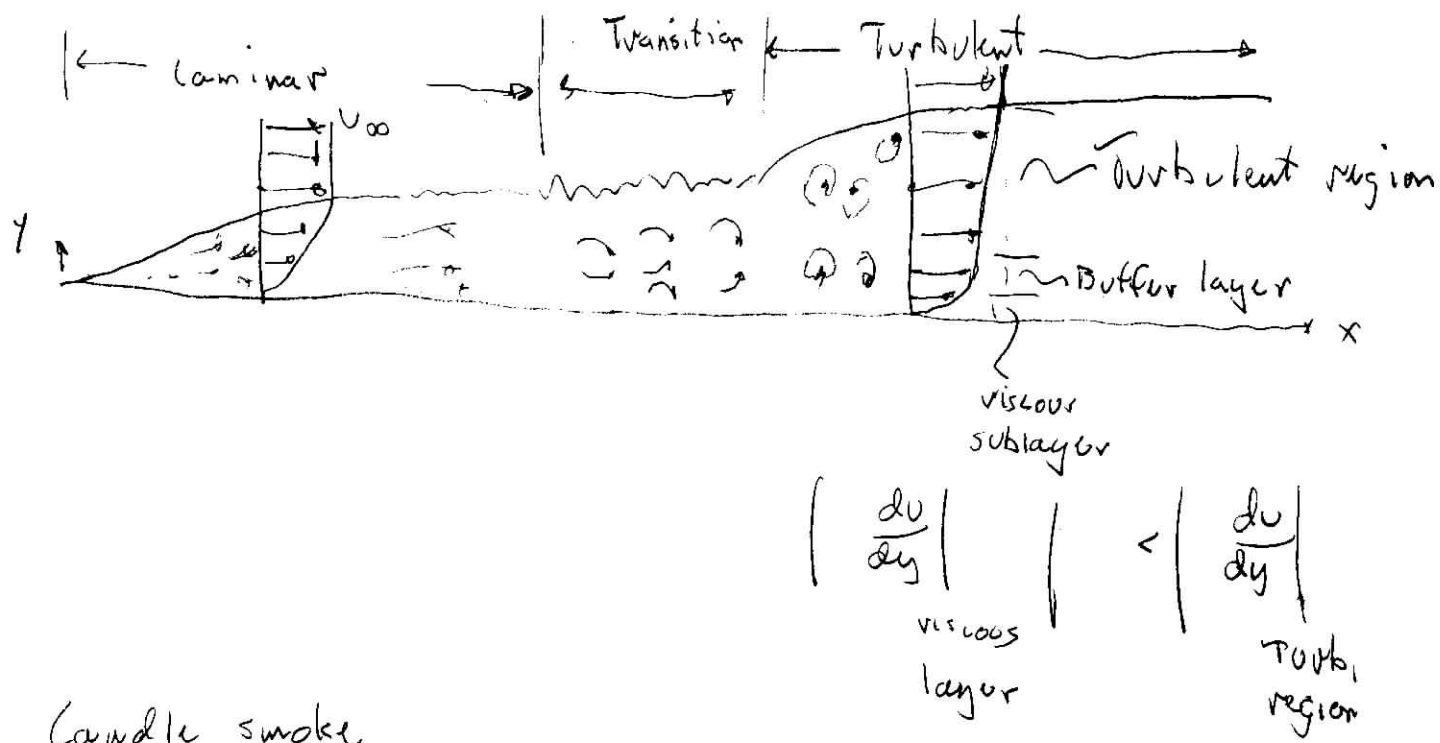
Water ~ 10

oils

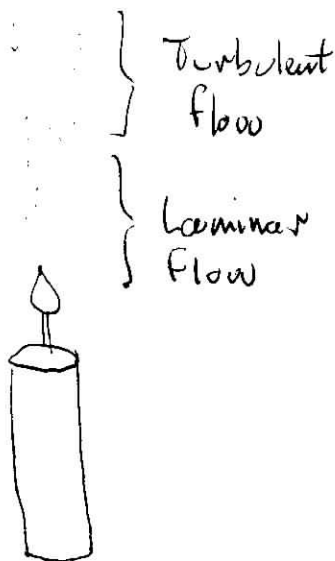
Small $\sim 10^{-3}$

Pr

länge $10^3 \sim 10^4$



Candle smoke



So what determines when
laminar flow $\rightarrow$ turb. flow

look at

$$Re \equiv \frac{\text{Inertial forces}}{\text{Viscous forces}} = \frac{U_\infty L_{char}}{\nu} = \frac{\rho U_\infty L_{char}}{\mu}$$

dimensionless

$[m^2/s]$

$$\sim \frac{\text{momentum dif.}}{\text{viscous dif.}}$$

$$\sim \frac{\text{overturning}}{\text{damping}}$$

Flat plate $Re_{crit} \sim 5 \times 10^5$

Pipe flow $\sim 2 \times 10^3$

$\sim$ varies with geometries & flow conditions

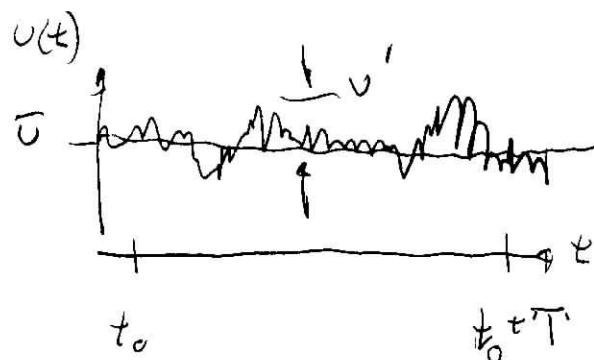
Average prop. value $\bar{U}$
 Inst. prop. value U
 flux. prop. value U'

$U = \bar{U} + U'$

inst. avg transient
 on flux. from $\bar{U}$

Average value of a property

$$\bar{U} \equiv \frac{1}{T} \int_{t_0}^{t_0+T} U(t) dt$$



Lots of calculations & definitions

... $\tau_{\text{turbulent}} = -\overline{\rho U' V'} = \mu + \frac{\partial \bar{U}}{\partial y}$

and so on.

$\dot{q}_{\text{turb}} = \rho C_p \sqrt{N T} = -k_{\text{turb}} \frac{\partial \bar{T}}{\partial y}$

lots of "second order" terms resulting from the $()'$ values

Ex) $\mu_{\text{turb}} \sim$ turbulent or eddy viscosity

So overall

$$\tau_{\text{total}} = (\mu + \mu_{\text{t}}) \frac{\partial \bar{U}}{\partial y} = \rho (\nu + \nu_{\text{t}}) \frac{\partial \bar{U}}{\partial y}$$

$$\dot{q}_{\text{total}} = -(k + k_{\text{t}}) \frac{\partial \bar{T}}{\partial y} = -\rho C_p (\alpha + \alpha_{\text{t}}) \frac{\partial \bar{T}}{\partial y}$$

Conservation Equations (Mass, Mom, Thermal)

$$\underline{V} = U(x,y)\hat{i} + V(x,y)\hat{j} + W(x,y)\hat{k}$$

Mass $\nabla \cdot \underline{V} = \frac{\partial U}{\partial x} + \frac{\partial W}{\partial y} = 0 \quad (\rho = \text{const})$

x-dir Mom. $\rho \frac{D\underline{V}}{Dt} = \rho \left(U \frac{\partial U}{\partial x} + V \frac{\partial W}{\partial y} \right) = \mu \frac{\partial^2 U}{\partial y^2} - \frac{\partial P}{\partial x} \quad (\text{const. } \rho, \mu)$

plus 2 more for y & z momentum comp.

$P = P(x)$ only for no grav. or other body forces

Thermal Energy $\rho C_p \left(U \frac{\partial T}{\partial x} + V \frac{\partial T}{\partial y} \right) = k \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + \mu \Phi$

$\Phi = \text{viscous dissipation}$

$$= 2 \left[\left(\frac{\partial U}{\partial x} \right)^2 + \left(\frac{\partial V}{\partial y} \right)^2 \right] + \left(\frac{\partial U}{\partial y} + \frac{\partial V}{\partial x} \right)^2$$

Homework

If viscous shear forces are negl. $\Phi \approx 0$ compared to

$$\rho C_p \left(U \frac{\partial T}{\partial x} + V \frac{\partial T}{\partial y} \right) = k \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right)$$

→ $\rho C_p \frac{DT}{Dt} = k \nabla^2 T$

$T = T(x,y)$

also

→ $\rho \frac{DU}{Dt} = \mu \nabla^2 U$

$U = U(y)$ only

→

$\nabla \cdot \underline{V} = 0$

$\underline{V} = U(y)\hat{i} + V(x,y)\hat{j} + 0\hat{k}$

Whole picture for a flat plate

farfield U_∞, T_∞

leading edge

U_∞, T_∞

y

x

plate

$$u(x, 0) = 0$$

$$v(x, 0) = 0$$

$$T(x, 0) = T_s$$

S_m or S_T

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2}$$

B.C.

① $x = 0$

leading edge
of plate

$$u(0, y) = U_\infty \quad T(0, y) = T_\infty$$

② $y = 0$

at plate

$$u(x, 0) = 0 \quad v(x, 0) = 0 \quad T(x, 0) = T_s$$

③ $y \rightarrow \infty$

far field

$$u(x, \infty) = U_\infty \quad T(x, \infty) = T_\infty$$

Note: that we never say $v(x, \infty) = 0$!!!

why?

We want

$$u = u(x, y; L, U_\infty, \rho, \mu,$$

$$T = T(x, y; L, U_\infty, \rho, \mu, T_s, T_\infty, k)$$

Non-dimensionalize

• characteristic length L_c or L
(plate length)

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• characteristic velocity U_∞

$$\tilde{x} = \frac{x}{L} \quad \tilde{y} = \frac{y}{L} \quad \tilde{U} = \frac{U}{U_\infty} \quad \tilde{N} = \frac{N}{U_\infty} \quad \tilde{P} = \frac{P}{\rho U_\infty^2} \quad \Theta = \frac{T(x, y) - T_s}{T_\infty - T_s}$$

Reduce to

Mass $\frac{\partial \tilde{U}}{\partial \tilde{x}} + \frac{\partial \tilde{N}}{\partial \tilde{y}} = 0$

x-Mom. $\tilde{U} \frac{\partial \tilde{U}}{\partial \tilde{x}} + \tilde{N} \frac{\partial \tilde{U}}{\partial \tilde{y}} = \frac{1}{Re_L} \frac{\partial^2 \tilde{U}}{\partial \tilde{y}^2} - \frac{d\tilde{P}}{d\tilde{x}}$

Thermal $\tilde{U} \frac{\partial \tilde{T}}{\partial \tilde{x}} + \tilde{N} \frac{\partial \tilde{T}}{\partial \tilde{y}} = \frac{1}{Re_L Pr} \frac{\partial^2 \tilde{T}}{\partial \tilde{y}^2}$

$$Re_L = \frac{\rho U_\infty L}{\mu}$$

$$Pr = \frac{\nu}{\alpha}$$

B.C.

$$\tilde{U}(0, \tilde{y}) = 1$$

$$\tilde{N}(\tilde{x}, 0) = 0$$

$$\tilde{T}(0, \tilde{y}) = 1$$

$$\tilde{U}(\tilde{x}, 0) = 0$$

$$\tilde{T}(\tilde{x}, 0) = 0$$

$$\tilde{U}(\tilde{x}, \infty) = 1$$

$$\tilde{T}(\tilde{x}, \infty) = 1$$

We are looking for

$$\tilde{U} = \tilde{U}(\tilde{x}, \tilde{y}; Re_L)$$

$$\tilde{T} = \tilde{T}(\tilde{x}, \tilde{y}; Re_L, Pr)$$

$\tilde{N}$ is an afterthought

At the end of the day we want things like

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$$\tau_{\text{wall}} = \mu \left. \frac{\partial u}{\partial y} \right|_{y=0} = \dots = \left(\frac{\mu U_{\infty}}{L} \right) \left. \frac{\partial \tilde{u}}{\partial \tilde{y}} \right|_{\tilde{y}=0} = \left(\frac{\mu U_{\infty}}{L} \right) f_1(\tilde{x}, Re)$$

$$C_{f,x} = \frac{\tau_{\text{wall}}}{\left(\frac{1}{2} \rho U_{\infty}^2 \right)} = \dots = \frac{(\mu U_{\infty}/L)}{\left(\frac{1}{2} \rho U_{\infty}^2 \right)} \left. \frac{\partial \tilde{u}}{\partial \tilde{y}} \right|_{\tilde{y}=0} = \frac{2}{Re} f_1(\tilde{x}, Re)$$

$$h_x = \frac{-k \left. \frac{\partial T}{\partial y} \right|_{y=0}}{(T_s - T_{\infty})} = \dots = \frac{k}{L} \left. \frac{\partial \tilde{T}}{\partial \tilde{y}} \right|_{\tilde{y}=0}$$

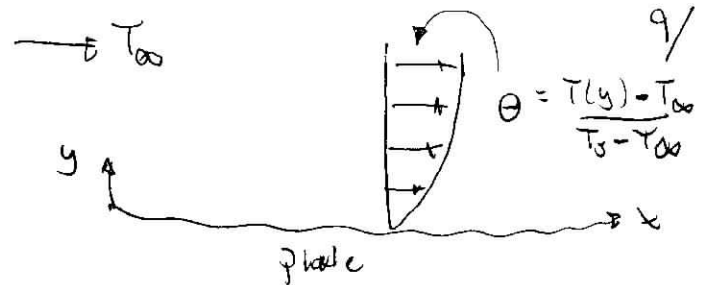
$$Nu_x = \frac{h_x L}{k} = \left. \frac{\partial \tilde{T}}{\partial \tilde{y}} \right|_{\tilde{y}=0} = f_2(\tilde{x}, Re, Pr)$$

Ex

Air @ $T_{\infty} = 20^{\circ}\text{C}$

flows over flat plate

@ $T_s = 160^{\circ}\text{C}$.



The dim. less temp is $\theta(y) = \frac{T(y) - T_{\infty}}{T_s - T_{\infty}} = e^{-ay}$

where $a = 3200 \text{ m}^{-1}$.

Determine - heat flux at plate surface

- convection heat transfer coef.

Soln Assume radiation is negl.

Calc. k air at average film temp $\bar{T}_{\text{air}} = \frac{T_s + T_{\infty}}{2} = 90^{\circ}\text{C}$

From tables $k = 0.03024 \frac{\text{W}}{\text{mK}}$
fluid

"
 q_{cond} to surface from plate interior
"
 q_{conv} away from plate by fluid

$$\begin{aligned} q_{\text{cond}} &= -k_f \left. \frac{\partial T}{\partial y} \right|_{y=0} = -k_f (T_s - T_{\infty}) a e^{-ay} \bigg|_{y=0} \\ &= (T_s - T_{\infty}) (-a k_f) \\ &= -4.48 \times 10^5 \frac{\text{W}}{\text{m}^2} (k_f) \\ &= 1.35 \times 10^4 \frac{\text{W}}{\text{m}^2} \end{aligned}$$

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Finally
$$h = \frac{-k_{\text{fluid}} \left(\frac{\partial T}{\partial y} \right)_{y=0}}{(T_s - T_\infty)} = 96.8 \frac{\text{W}}{\text{m}^2 \text{K}}$$

Note one could also use Newton's law of cooling

$$q'' = h(T_s - T_\infty)$$

